

(1)

### Introduction

Our project pertains to a quantum system where at time t = 0 we know that with probability one the information is at node l. We track the time evolution of the system in order to find a time T where the information is at node m with probability 1 if there exists such a time T. We wish to characterize matrices that realize Perfect State Transfer. Specifically we are looking at matrices that represent cycles on nnodes. To categorize these matrices we have examined the spectrum of such matrices. Additionally we have proved that an algorithm that splits paths on n vertices into cycles on 2n-2 preserves Perfect State Transfer.

#### Definitions

1. Consider

$$\mathcal{K} = \begin{bmatrix} 0 & k_1 & 0 & \dots & 0 & x \\ k_1 & 0 & k_2 & 0 & \dots & 0 \\ 0 & k_2 & \ddots & \ddots & \ddots & \ddots & 0 \\ \vdots & 0 & \ddots & \ddots & \ddots & 0 \\ 0 & \vdots & \ddots & 0 & k_{n-1} \\ x & 0 & \dots & 0 & k_{n-1} & 0 \end{bmatrix}$$

This form represents the adjacency matrix of a cycle graph with weights. If x = 0 we have a matrix representation of a path on n nodes.

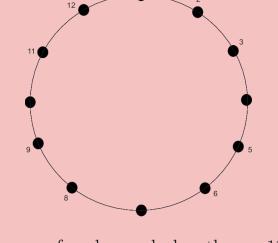


Image of cycle graph, length n = 12

#### Note $\mathcal{K}$ is called a **Periodic Jacobi Matrix**.

2. **Perfect State Transfer**(PST) occurs between nodes l and m if and only if there exists a time t such that

$$e^{itK}e_{\ell} = \varphi e_m$$

where  $\varphi \in \text{and } |\varphi| = 1$ .

#### Theorem 1

Let  $\mathcal{K}$  be an  $n \times n$  symmetric matrix with real entries and let  $\vec{v_k} = (v_{k,1}, v_{k,2}, \dots, v_{k,n})$ be the k-th eigenvector of  $\mathcal{K}$ ,  $1 \leq k \leq n$ , with corresponding eigenvalue  $\lambda_k$ . The following are necessary and sufficient for Perfect State Transfer to occur between nodes  $\ell$  and m.

1.  $|\vec{v}_{k,\ell}| = |\vec{v}_{k,m}|$  for all k and

2. 
$$\lambda_k - \lambda_{k+1} = \frac{h_k \pi}{T}$$
, with

• 
$$h_k \in 2\mathbb{Z}$$
 if  $\frac{\vec{v}_{k,l}}{\vec{v}_{k,m}} = \frac{\vec{v}_{k+1,l}}{\vec{v}_{k+1,m}}$ 

- $h_k \in 2\mathbb{Z} + 1$  if  $\frac{\vec{v}_{k,l}}{\vec{v}_{k,m}} = -\frac{\vec{v}_{k+1,l}}{\vec{v}_{k+1,m}}$

### Theorem 2

Let  $\mathcal{K}$  be as in equation (1). If  $\mathcal{K}$  does not have a simple spectrum, then  $\mathcal{K}$  does not realize perfect state transfer from node 1 to node n.

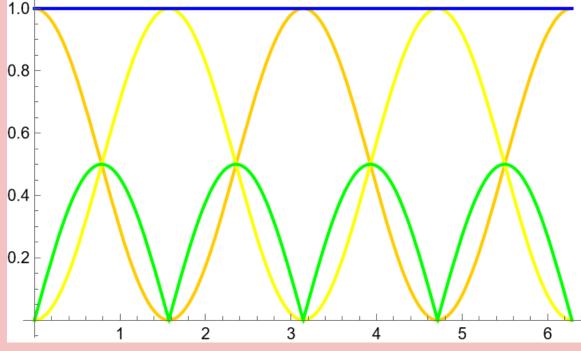
# QUANTUM PERFECT STATE TRANSFER IN CYCLE GRAPHS Elizabeth Athaide<sup>1</sup>, Leia Donaway<sup>2</sup>, Sam Trombone<sup>3</sup> <sup>1</sup>Massachusetts Institute of Technology, <sup>2</sup>Swarthmore College, <sup>3</sup>Hamilton College

Cons

$$\begin{pmatrix} x & 0 & 1 & 0 \end{pmatrix}$$

**Example**  

$$\begin{aligned} & \text{Theorem 3} \\ \text{the states integers  $h_{2i} = h_{2i} + h$$$



1. Let 
$$\mathcal{P} = \begin{pmatrix} 0\\b_1\\0\\0 \end{pmatrix}$$

**Example**  
addet the matrix 
$$C = \begin{pmatrix} 0 & 1 & 0 & 2 \\ 1 & 1 & 1 & 0 \\ 2 & 1 & 1 & 0 \\ 1 & 1 & 1 & 0 \\ 2 & 1 & 1 & 0 \end{pmatrix}$$
.  
The differences of vigne values  $h_1 = h_1 = h_1$ ,  $h_2 = h_1 = h_2 = h_1 + h_2 = h_1 + h_2 = h_2 + h_2$ 

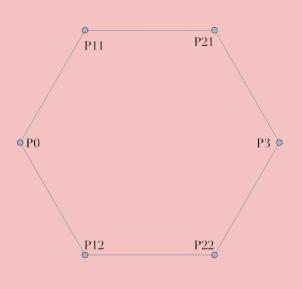
- 3. Each  $P_i$  corresponds to a node on  $P_4$ .

We will break this path into a cycle by splitting each inner node,  $P_i$ , into two nodes,  $P_i^1, P_i^2$ , where,

$$P_i^1 + P_i^2 = P_i$$
 and  $P_i^1 = P_i^2$ .

4. Our new cycle can be visualized as





We hope to characterize periodic Jacobi matrices that realize perfect state transfer.

- Is it always the case that a Hamiltonian that realizes perfect state transfer on a cycle can be reduced to a Jacobi matrix?
- Is our construction the only Hamiltonian that realizes perfect state transfer between nodes 1 and n in a cycle of length 2n - 2?

## **Acknowledgements and References**

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